INVESTIGATION OF THE BEHAVIOUR OF UNSTEADY TWO-DIMENSIONAL PERIODIC FLOWS DEFINED BY THE EQUATIONS OF MOTION

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As a mathematical model for approximate transfer of gas motions from the full system of Navier-Stokes equations, only the equations of motion are studied assuming constant values of thermodynamic parameters of density and temperature: $\rho = 1$, T = 1.

The velocity scale u_{00} is taken as $1/3 \cdot 10^3$ M/c, which is close to the speed of sound in air under normal conditions. The distance scale x_{00} is taken as the value corresponding to the geometric characteristics of the particular flow under study.

This paper considers the case of absence of z dependence and equality to zero of the third component of the gas velocity vector: $\partial/\partial z = 0$; $\upsilon_3 = 0$.

In this case, the system under consideration in a detailed form has the following form

$$\begin{cases} u_{t} = -uu_{x} - \upsilon u_{y} + \mu_{0} \left(u_{xx} + \frac{3}{4} u_{yy} + \frac{1}{4} u_{xy} \right), \\ \upsilon_{t} = -u\upsilon_{x} + \upsilon\upsilon_{y} + \mu_{0} \left(\upsilon_{yy} + \frac{3}{4} \upsilon_{xx} + \frac{1}{4} u_{xy} \right). \end{cases}$$
(1)

Further the system (1) will be referred to as a system of equations of motion. In this paper, the case of two spatial variables is considered and, taking into account the results from [1-5], the following representations of the sought functions are used u, v.

$$u(t, x, y) = \sum_{k=1}^{K} \left[u_{k,1} \sin(kx) + u_{k,2} \sin(ky) \right],$$

$$\upsilon(t, x, y) = \sum_{k=1}^{K} \left[\upsilon_{k,1} \sin(kx) + \upsilon_{k,2} \sin(ky) \right].$$
(2)

With the help of constants $u_{k,1}^0$, $u_{k,2}^0$, $v_{k,1}^0$, $v_{k,2}^0$ at time t = 1 initial conditions are set for systems of ordinary differential equations in which the upper indices in the sums are taken equal to K, i. e., M = K. The variable ℓ in these systems takes integer values from one to K.

$$u_{\ell,1}' = -\frac{1}{2} \sum_{k=1}^{\infty} \sum_{m=1}^{\infty} m u_{k,1}(t) u_{m,1}(t) b_{k,\ell,m} - \mu_0 \ell^2 u_{\ell,1}; \qquad \ell = 1, 2, 3, \dots$$
(3)

$$\upsilon_{\ell,1}' = -\frac{1}{2} \sum_{k=1}^{\infty} \sum_{m=1}^{\infty} m u_{k,1}(t) \upsilon_{m,1}(t) b_{k,\ell,m} - \frac{3}{4} \mu_0 \ell^2 \upsilon_{\ell,1}; \quad \ell = 1, 2, 3, \dots$$
(4)

$$u_{\ell,2}' = -\frac{1}{2} \sum_{k=1}^{\infty} \sum_{m=1}^{\infty} k \upsilon_{m,2}(t) u_{k,2}(t) b_{\ell,m,k} - \frac{3}{4} \mu_0 \ell^2 u_{\ell,2}; \quad \ell = 1, 2, 3, \dots$$
(5)

$$\upsilon_{\ell,2}' = -\frac{1}{2} \sum_{k=1}^{\infty} \sum_{m=1}^{\infty} k \upsilon_{k,2}(t) \upsilon_{m,2}(t) b_{\ell,m,k} - \mu_0 \ell^2 \upsilon_{\ell,2}; \quad \ell = 1, 2, 3, \dots$$
(6)

The obtained Cauchy problems for the systems are solved numerically at $0 \le t \le t_f$, where the finite moment t_f is chosen based on the meaning of the problem under consideration and its solutions.

References

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